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DFT Study of Type-II SnBr₂/Cul vdW Heterostructure: Implications for Thermophotovoltaic Devices

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ABSTRACT: In present investigation, first-principle investigation on the structural stability and tunable optoelectronic characteristics of a SnBr₂/CuI van der Waals heterostructure (vdW HTS). The stability of the vdW HTS is also verified by adhesion energy calculations and ab initio molecular dynamics (AIMD) simulations, exhibiting its energetic and thermal stability. Electronic band structure calculations determine SnBr₂/CuI vdW HTS to exhibit a direct bandgap near 1.17 eV within the HSE06 functional, in the optimum bandgap regime for solar absorbers. Crucially, the type-II band alignment enables spatial separation of photogenerated holes and electrons, thus inhibiting recombination and favoring carrier transportation across the interface. Charge density difference analysis also validates interfacial charge reconstructions in accordance with efficient carrier exchange. The optical calculations indicate a clear red-shifted absorption edge into the visible regime, along with an increased absorption coefficient over the pure monolayer counterparts, implying enhanced light-harvesting capability. Together, these findings establish the SnBr₂/CuI vdW HTS to offer excellent structural stability alongside tunable optoelectronic performance, thus making it an extremely promising next-generation solar and nano-optoelectronic contender.

Keywords: vdW Heterostructure, Thermophotovoltaic, DFT, Type-II semiconductor.

INTRODUCTION

To achieve efficient carrier separation in photoelectric, photovoltaic devices, it is crucial to explore suitable materials and structures. One promising solution is the use of van der Waals heterostructure (vdw HTS) with a type-II band alignment. These HTSs consist of twodifferent monolayers (MLs) stacked on top of each other without forming chemical bonds in between. The concept of vdw HTS was first introduced by Geim et al. 2013, and it has gained significant attention in the field of materials research (Meng et al., 2021). The weak vdw forces between these monolayers allow them to retain their original monolayer's optoelectronic characteristics, while the interlayer interaction brings out the novel characteristics. Many researchers have successfully constructed various 2D vdw HTS such as PtS2/ZrS2 (Parmar et al., 2023), GaTe/ZnI₂, PtSe₂/GaSe (Parmar et al., 2023), and many others. The discovery of graphene (Novoselov et al., 2005) has brought new insight into 2D materials. After that, many 2D materials are studied extensively because of their novel properties and potential applications (Khengar et al., 2022, Modi et al., 2023). Of all these various 2D materials, halidesdihalides have gained research interest due to their interesting properties, potential applications, and tunable bandgap (Hu et al., 2020). Recently, SnBr2 has been theoretically investigated, having an indirect bandgap of 3.02 eV (Naik et al., 2023). While CuI has a direct bandgap having an application in photovoltaic devices (Amrani et al., 2006). Due to these excellent properties

of SnBr₂ and CuI MLs, we choose to construct SnBr₂/CuI vdw HTS. Recently, 2D materials are explored for many novel applications such as Gas sensors (Peng et al., 2021), storage devices (Mortazavi et al., 2016), Thermophotovoltaic (TPV) devices (Ismail et al., 2024), photocatalysis (Shakil et al., 2023), etc. Among which, TPV is gaining more research interest due its various potential use. TPV devices are energy conversion systems that directly transform thermal radiation into electricity which is the same fundamental principle as photovoltaic cells. Such TPV devices not only use sunlight just like solar cell but it harnesses the thermal radiation emitted from high-temperature sources such as combustion systems, concentrated solar absorbers, nuclear reactors, or even industrial waste heat. A typical TPV system consists of a thermal emitter, which radiates a broad spectrum of photons when heated to elevated temperatures, and a photovoltaic cell that absorbs these photons to generate electron-hole pairs and produce electrical power (Agarwal and Prajapati 2018). Spectral control elements such as filters or photonic structures are used in TPV which reflects subbandgap photons back to the emitter and hence efficiency improved.

The efficiency of TPV devices strongly depends on the choice of semiconductor material for the photovoltaic cell, as the bandgap must be carefully matched with the spectral distribution of the thermal emitter (Ahmed *et al.*, 2025). Low bandgap semiconductors in the range of 0.2–0.7 eV, including materials such as InGaAsP (Tuley and Nicholas 2010), GeTe (Abir *et al.*, 2023), have been

widely studied because of their ability to absorb nearinfrared thermal radiation effectively. In addition, achieving direct bandgaps and favorable carrier dynamics through Type-II band alignment in heterostructures has been shown to improve photon absorption and charge extraction, both of which are essential for efficient device performance (Almayyali et al., 2023). In this context, recent advances in twodimensional (2D) materials and van der Waals heterostructures have opened new possibilities for TPV device design, owing to their tunable band structures, strong light-matter interactions, and potential to achieve Type-II band alignment with direct bandgaps. These features make them promising candidates for nextgeneration TPV applications.

Currently, no systematic research has been found on the SnBr₂/CuI vdW Heterostructure for the implications in Thermophotovoltaic Devices. So, this research consists of the DFT investigation on SnBr₂/CuI vdw HTS thermophotovoltaic Devices.

COMPUTATIONAL DETAILS

In this research, all the simulations have been performed within the DFT framework (Orio *et al.*, 2009) by Quantum Espresso software. The exchange-correlation functional of the electron is described by utilizing the Generalized Gradient Approximation (GGA) given by the Perdew-Burke-Eruzerhof (PBE) (Perdew *et al.*, 1996). To obtain a more accurate bandgap, a hybrid

functional HSE06 given by Heyd-Scuseria-Ernzerhof (Heyd *et al.*, 2003) is utilized. The DFT-D3 (Grimme *et al.*, 2011) method given by Grimme has been utilized to represent the vdw interactions. A thick vacuum of 25 Å is given in a perpendicular direction in order to prevent interactions between two layers. The energy and force threshold is set to 10^{-8} Ry and 10^{-4} Ry/Bohr, respectively, to obtain structural optimization and system convergence. The Monkhorst-Pack k-point grid is fixed at $6\times6\times1$ and plane wave energy cut-off is fixed at 550 Ry. The optical properties are obtained by using the Norm-conserving pseudopotential and frequency-dependent complex dielectric function ϵ (ω).

RESULTS AND DISCUSSIONS

A. Structural Properties

The monolayers $SnBr_2$ and CuI have similar hexagonal symmetry and almost the same lattice parameter, so it feasible to construct heterostructure easily. The obtained lattice constant (a) of $SnBr_2/CuI$ vdw HTS, $SnBr_2$ ML, and CuI ML are 4.32 Å, 4.37 Å, and 4.33 Å, respectively. The lattice mismatch between their constitute monolayer for $SnBr_2/CuI$ vdw HTS is given by equation $\delta = \frac{a_{CuI} - a_{SnBr_2}}{a_{CuI}}$ (Parmar *et al.*, 2023), which is 1.14% for this vdw HTS. The optimized structure of constructed $SnBr_2/CuI$ vdw HTS is illustrated in Fig. 1(a).

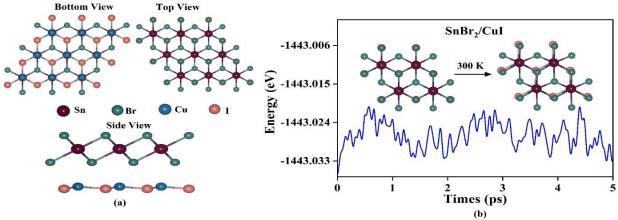


Fig. 1. (a) Optimized structure and (b) AIMD calculation of SnBr₂/CuI vdw HTS.

To check the energetic stability of this constructed vdw HTS, we have to find its adhesion energy E_{ae} that is given by equation (Parmar *et al.*, 2022), $E_{ae} = E_{\nu dw\ HTS} - E_{SnBr_2} - E_{Cul}$. Where $E_{vdw\ HTS}$ is the total energy of a HTS, and E_{SnBr_2} and E_{Cul} represent the energy of their respective monolayers. The obtained negative value of adhesion energy -1.01 eV for the SnBr₂/CuI vdw HTS shows that the material is energetically stable. However, the AIMD simulation is also performed to ensure the thermal stability of SnBr₂/CuI vdw HTS for 5 ps with a time interval of 1fs. The obtained results as shown in Fig. 1(b), illustrates that there is minimal deviation in total energy and negligible structural fluctuations throughout the entire simulation. Hence, the constructed SnBr₂/CuI vdw HTS is energetically and thermally stable material.

It can also be added that the extremely low lattice mismatch (1.14%) implies negligible interfacial strain, favorable for the retention of the intrinsic quality of the constituting layers while allowing for tight interlayer coupling (Guo et al., 2017). The geometrically optimized structure implies that the van der Waals interaction maintains the interlayer separation without structural distortions, confirming structural soundness. Such low values for the mismatch also favour easier experimental realization, due to the lack of strain-related defects. Furthermore, the stability also agreed by the roomtemperature AIMD implies the potential for the heterostructure to survive real-device operation conditions. Collectively, these results imply that the SnBr₂/CuI vdW HTS possesses the dual characteristic of robustness and compatibility, preconditions for the

realization of practical optoelectronic as well as photovoltaic device fabrication.

B. Electronic Properties

Since SnBr₂ ML, CuI ML, and SnBr₂/CuI vdw HTS have hexagonal symmetry, the band diagram is obtained along highly symmetric direction 'Γ-M-K-Γ' with HSE06 functional as shown in Fig. 2(a,b,c). The SnBr₂ ML has an indirect bandgap of 3.02 eV, while CuI ML has a direct bandgap of 2.50 eV. Interestingly, SnBr₂/CuI vdw HTS exhibits a direct bandgap of 1.21 eV, which is lesser than each monolayer of constructive components. The CBM resembles SnBr₂ while VBM mainly originates from CuI. Hence, in the presence of radiation, the SnBr₂/CuI vdw HTS exhibits efficient electron transfer, resulting in the effective separation of photoexcited electron-hole pairs. These findings also highlight the enhanced optical properties. The Fig. 2(d) displays the Projected Density of State (PDOS).

It is seen that Cu-d state is more dominant in valance band while Br-p and Sn-p states have significant effect in conduction band. Additionally, the overlap of Cu-d states with Sn-p/Br-p states around the Fermi level indicates interlayer hybridization, which enhances interfacial coupling and supports efficient charge transfer. The band curvature in Fig. 2 also suggests relatively light effective masses for electrons in the SnBr₂ conduction band and holes in the CuI valence band. Lower effective masses are typically associated

with higher carrier mobilities, which is a desirable feature for fast charge transport in photovoltaic applications. The materials work function Φ is characteristics as the energy required to eliminate an electron from its surface is obtained by employing the equation, $\Phi = E_{vac} - E_f$, where E_f and E_{vac} denote the fermi and vacuum level. The obtained Φ is 8.71 eV for vdw HTS.

As shown in Fig. 3(a), the CuI ML exhibits a higher work function compared to the SnBr₂ ML, which indicates that upon contact electrons will transfer from CuI to SnBr₂ surface until their fermi level matches. Furthermore, the charge density difference $\Delta \rho$ is obtained (Parmar *et al.*, 2023) from $\Delta \rho = \Delta \rho_{vdw \, HTS}$ – $\Delta \rho_{SnBr_2} - \Delta \rho_{CuI}$. The yellow clouds below SnBr₂ ML indicate charge accumulation, while green spheres above CuI ML suggest charge depletion, implying substantial charge transfer from CuI ML to SnBr₂ ML, as depicted in Fig. 3(b). The redistribution of charges forms an intrinsic inbuilt electric field at the junction, which in turn drives the photogenerated carriers away and hinders the loss due to recombination. The generated dipole layer at the interface at the end adjusts the electrostatic potential at the junction, favoring enhanced type-II band alignment and easier carrier separation. Such effects demonstrate the working principle of a normal p-n junction but in the pure 2D vdW regime (Chapin et al., 1954).

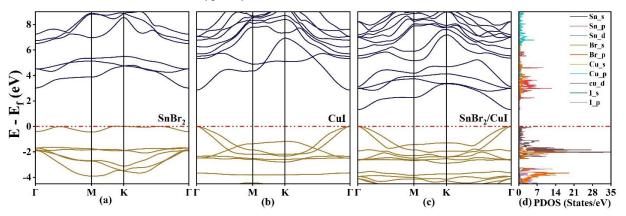


Fig. 2. Band structure of (a) SnBr2 ML (b) CuI ML (c) SnBr2/CuI vdw HTS and (d) PDOS of SnBr2/CuI vdw HTS.

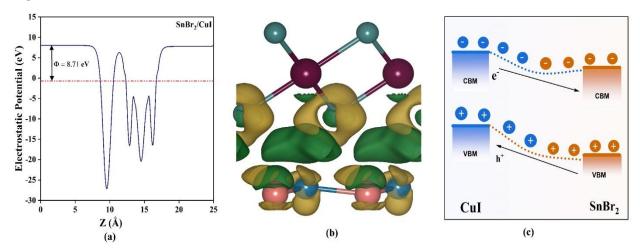


Fig. 3. (a) Work Function (b) Charge density (c) Charge transfer with band alignment.

Besides the photovoltaic uses, the offset alignment can be utilized for photocatalytic reaction, where spatially separated carrier pairs are utilized to drive reduction and oxidation processes on different layers (Dalsaniya *et al.*, 2021).

The band alignment is displayed in Fig. 3(c). Under solar light, the SnBr₂/CuI vdw HTS exhibits photogenerated electron transfer from CuI ML to the SnBr₂ conduction band and photoexcited holes move to CuI valance band, achieving spatial electron-hole separation. This is attributed to the type-II band alignment, effectively reducing electron-hole recombination.

C. Optical Properties

The optical response is calculated from the dielectric function (that is the function of frequency) $\varepsilon(\omega)$ given as $\varepsilon(\omega) = \varepsilon^1(\omega) + i\varepsilon^2(\omega)$. Here, $\varepsilon^1(\omega)$ and $\varepsilon^2(\omega)$ corresponds to the real and imaginary portion, respectively. The real part in Fig. 4(a) shows that the static dielectric constant (SDC) of SnBr₂ ML, CuI ML, and SnBr₂/CuI vdw HTS is 2.54, 1.76, 2.91, respectively. The increment in SDC of $\varepsilon^1(\omega)$ implies that lower exciton binding energy which leads to an effective

hole and electrons separation. The presence of a negative value indicates the metallic characteristics at that particular frequency. The dielectric function's imaginary portion is related to the interband transition, which is shown in Fig. 4(b) exhibits several distinct peaks. Further, the relation between absorption coefficient and energy is displayed in Fig. 4(c). and is obtained the equation (Mehta *et al.*, 2023),

$$\alpha(\omega) = \left(\frac{\sqrt{2}\omega}{c}\right) \left[(\varepsilon^1(\omega)^2 + \varepsilon^2(\omega)^2)^{\frac{1}{2}} - \varepsilon^1(\omega) \right]^{\frac{1}{2}}.$$

The highest peak is observed at 6.74 eV, 5.54 eV, and 7.55 eV for SnBr₂ ML, CuI ML, and SnBr₂/CuI vdw HTS, respectively. It is seen from the figure that the absorption of SnBr₂/CuI vdw HTS is enhanced from CuI ML. Also, SnBr₂/CuI vdw HTS exhibits a starting noticeable absorption in visible region in comparison with SnBr₂ ML. So, by comparing the absorption curve of both MLs with SnBr₂/CuI vdw HTS, it is seen that SnBr₂/CuI vdw HTS has better absorption performance. So, SnBr₂/CuI vdw HTS has great potential for efficient optoelectronic-nanoelectronics devices and photovoltaic application.

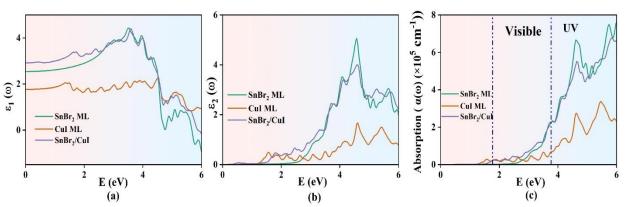


Fig. 4. (a) Real portion (b) Imaginary portion and (c) Absorption coefficient.

Besides the improved absorption, the red-shifted absorption edge for SnBr2/CuI vdW HTS holds special significance, for it enables the material to exploit a larger fraction of the solar spectrum than either of the individual layers. The relatively larger dielectric constant also translates to lower exciton binding energy, implying that the excitons can be readily split into free carriers—a characteristic important for photovoltaic conversion. Lastly, the sharp absorption peaks in the UV regime point to possible suitability for UV photodetectors (Patel et al., 2021). The simultaneous occurrence of visiblerange absorption and UV sensitivity also serves to reinforce the optoelectronic tunability of the heterostructure, making it suitable for various optoelectronic applications (Yang et al., 2019). Such tunability is characteristic for spectral vdW heterostructures and fortifies the candidacy for multifunctional devices for SnBr2/CuI.

CONCLUSIONS

In this work, we have systematically investigated the structural and optoelectronic properties of the SnBr₂/CuI van der Waals heterostructure (vdW HTS) using first-principles density functional theory calculations. Our results confirm that the heterostructure is both

energetically and thermally stable, as evidenced by the negative adhesion energy and the minimal fluctuations observed in AIMD simulations. The compatibility of lattice constants between SnBr₂ and CuI enables the formation of a stable interface with minimal mismatch, further supporting its structural integrity.

From an electronic point of view, SnBr₂/CuI vdW HTS has a direct bandgap in the region of about 1.17 eV, within the optimum value for photovoltaic absorbers. Note well that the type-II band alignment for efficient charge transfer over the interface takes place due to the spatial separation of electrons and holes in multiple layers. The band alignment strongly suppresses the recombination losses, thus increasing the probable efficiency in devices by virtue of this heterostructure. Charge density difference calculations and work functions are additional proof to confirm the taking place of the interfacial charge redistribution in support of effective carrier separation.

Optical characterizations reveal weak absorption in the visible regime for the pure SnBr₂ monolayer, while the SnBr₂/CuI heterostructure red-shifts the absorption edge to align the absorption within the visible regime. The above-mentioned features along with the relatively low value of the exciton binding energy indicated by the

larger dielectric constant direct towards efficient separation and generation of the photoexcited carrier pairs in the heterostructure. Consequently, the synergistic combination of inherent structural stability, appropriate band alignment, and improved optical response makes the SnBr₂/CuI vdW HTS an extremely viable prospect for next-generation photovoltaic and nano-optoelectronic devices. The SnBr₂/CuI vdW heterostructure demonstrates type-II band alignment and a direct bandgap, both of which are highly favorable for efficient carrier separation and strong light absorption. These features suggest that the system could be a thermophotovoltaic promising candidate for applications, where optimized bandgaps and charge dynamics are essential for high conversion efficiency.

FUTURE SCOPE

The SnBr2/CuI heterostructure from the present study has the appropriate stability, band alignment, and optical response for next-generation photovoltaic and TPV devices. Further simulations at the device level, experimental synthesis, and tuning by strain, doping, or external fields are suggested to further optimize the performance of SnBr2/CuI. Its strong interlayer charge separation and visible-range absorption also make it a good candidate for broader optoelectronic applications.

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Conflict of Interest. The authors declare that they have no known financial or personal relationships that could have influenced the work reported in this paper. No competing interests exist.

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